

# Laser induced elastico mechano luminescence of $\text{SrAl}_2\text{O}_4 : \text{Eu}$ phosphor

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## Abstract

When  $\gamma$  – irradiated elastic mechano luminescent materials  $\text{SrAl}_2\text{O}_4 : \text{Eu}$  are exposed to 1060 nm infrared pulse of nanosecond duration from  $\text{CO}_2$  laser, then stress produced in the crystals excites visible luminescence due to piezoelectric field in  $\text{SrAl}_2\text{O}_4 : \text{Eu}$  phosphor because they are non – centro symmetric. In the present investigation  $\text{SrAl}_2\text{O}_4 : \text{Eu}$  crystals are given laser shocks and ML intensity recorded. During laser induced shocks, ML intensity increases linearly with stress and attains a peak value at a particular time and then decays exponentially with time. A theoretical approach has been proposed to explain the experimental results.

**Keywords:** Elastico ML, phosphor, infrared laser.

## INTRODUCTION

Mechanoluminescence (ML) can be excited by impulsive deformation of solids [1]. The ML can be excited by the shock – waves produced during exposure of the materials to laser pulses [5]. When crystals are cleaved by laser, light emission takes place known as mechanoluminescence (ML) [1-4]. On cleavage of crystals, charged surfaces are produced which may cause dielectric breakdown of the solids and the intervening gases. Thus, light emission may take place. In the case of  $\text{SrAl}_2\text{O}_4 : \text{Eu}$  crystals, the ML measurements after high purification has not been made, but it has been found that for lower concentration of  $\text{Eu}^{2+}$  ions in  $\text{SrAl}_2\text{O}_4$ , the ML intensity increases drastically with increasing concentration of  $\text{Eu}^{2+}$ .

## THEORY

In the measurement of the elastico ML of  $\text{SrAl}_2\text{O}_4$  phosphors, EML is induced by laser light for pressing the samples. If  $\Omega$  is the activation volume,  $n_1$  is the number of effective defect centres in a crystal,  $n_c$  is the number of crystallites in the sample, and  $n_t$  is the concentration of filled electron traps, then the total number of detrapable traps is  $n_{t0} = \Omega n_1 n_t n_c$ . If  $d_0$  is the local piezoelectric constant near the effective defect centres, then for the applied pressure  $p$ , the piezoelectric charge  $q$  near the effective defect centres is given by,  $q = d_0 p$ . If the crystal is compressed at a fixed pressing rate  $\dot{p}$  or strain rate  $\dot{\epsilon}$ , then in the elastic region,  $q$  is given by

$$q = d_0 p = d_0 \dot{p} t = d_0 Y \dot{\epsilon} t \quad \dots\dots(1)$$

where  $Y$  is the Young's modulus of the elasticity of the crystal. If  $f_c$  is the characteristic piezoelectric field then we can write the following equation

$$-\frac{dn_t}{df} = \frac{n_t}{f_c} - zn_t \quad \dots\dots(2)$$

where  $f$  is the piezoelectric field near  $\text{Eu}^{2+}$  ions,  $n_t$  is the number of filled electron traps at any time  $t$ , and  $z = 1/f_c$ .

Integrating Eq. (2) and taking  $n_t = n_{t0}$  for the threshold field  $f = f_{th}$  we get

$$n_t = n_{t0} \exp [-z(f - f_{th})] \quad \dots\dots(3)$$

Where  $n_{t0}$  is the total number of the filled electron traps in the activation volume  $\Omega n_1$ .

Using Eq. (3), the total number of detrapped electrons can be expressed as

$$n_d = (n_{t0} - n_t) = n_{t0} [1 - \exp\{-z(f - f_{th})\}] \quad \dots\dots(4)$$

In the elastic region,  $z(f - f_{th})$  is much less than 1, hence, Eq. (4) can be written as

$$n_d = n_{t0} [1 - 1 + z(f - f_{th})] = n_{t0} z(f - f_{th}) \quad \dots\dots(5)$$

Now, differentiating Eq. (5), we get

$$\frac{dn_d}{dt} = n_{t0} z \frac{df}{dt} \quad \dots\dots(6)$$

As the rate of generation  $g$  of electrons in the conduction band will be equal to the rate of detrapping of electrons, we can write

$$g = n_{t0} z \frac{df}{dt} \quad \dots\dots(7)$$

If  $\tau$  is the lifetime of electrons in the conduction band, then the change in the number of electrons in the conduction band can be expressed as

$$\Delta n = g\tau = n_{t0} z \frac{df}{dt} \tau \quad \dots\dots(8)$$

Using Eq. (8), the current density  $j$  flowing in the crystal can be written as

$$j = \Delta n q v_d \quad \dots\dots(9)$$

where  $q$  is the electronic charge and  $V_d$  is the drift velocity.

If  $\mu$  is the mobility of electrons in the crystals, then  $V_d = \mu f$ , and Eq. (9) can be expressed as

$$j = \Delta n q \mu f \quad \dots(10)$$

Thus, the rate  $r$  for the flow of electrons in the conduction band in the crystal is given by

$$r = \frac{j}{q} = \Delta n v_d = \Delta n \mu f \quad \dots(11)$$

From Eqs. (8) and (11), we get

$$r = \Delta n v_d = n_{t0} z \tau \mu f \frac{df}{dt} \quad \dots(12)$$

If  $\sigma$  is the capture-cross section of the energy state for the excited  $\text{Eu}^{2+}$  ions, and  $n_c$  is the concentration of  $\text{Eu}^{2+}$  centres, the rate of generation of excited  $\text{Eu}^{2+}$  ions can be expressed as

$$R = \sigma n_c \Delta n v_d = \sigma n_c n_{t0} z \tau \mu f \frac{df}{dt} \quad \dots(13)$$

If  $\eta$  is the efficiency for the radiative decay of excited  $\text{Eu}^{2+}$  ions, then the ML intensity can be expressed as

$$\begin{aligned} I &= \eta R = \eta \sigma n_c n_{t0} z \tau \mu (f - f_{th}) \frac{df}{dt} \\ \text{or, } I &= \eta \sigma \Omega n_1 n'_c n_t n_c z \tau \mu (f - f_{th}) \frac{df}{dt} \end{aligned} \quad \dots(14)$$

Where  $f_{th}$  is the threshold piezoelectric field for the ML emission.

It is to be noted that  $n_c$  is related with  $n_1$ , hence, there should be nonlinear relation between  $I$  and  $n_c$ . A nonlinear relation between  $I$  and  $n_c$  has been observed. If  $b$  is the co-relating factor between the piezoelectric field  $f$  and the piezoelectric charge  $q$ , then,  $f = bq$ , and Eq. (14) can be written as

$$I = \eta \sigma \Omega n_1 n'_c n_t n_c z \tau \mu b^2 (q - q_{th}) \frac{dq}{dt} \quad \dots(15)$$

where  $f_{th} = b q_{th}$

Now, we will consider the EML Under Increasing Stress and Fixed Stress Condition- Case I – Rise of EML intensity. From Eqs. (1) and (15), we get

$$I = \eta \sigma \Omega n_1 n'_c n_t n_c z \tau \mu b^2 d_0^2 (p - p_{th}) \frac{dp}{dt} \quad \dots(16)$$

In terms of strain rate, Eq. (16) can be expressed as

$$I = \eta \sigma \Omega n_1 n'_c n_t n_c z \tau \mu b^2 d_0^2 Y \dot{\epsilon} (p - p_{th}) \quad \dots(17)$$

Equation (17) indicates that for a given strain rate the EML intensity should increase linearly with the pressure, and for a given pressure, the EML intensity should increase linearly with the strain rate.

### Case – II Decay of EML intensity

When the laser impulse has been stopped then decay of EML is related to the decrease of the strain rate of the sample with time. If

$\tau_m$  is the time-constant then the decrease of strain rate with time can be expressed as

$$\dot{\epsilon} = \dot{\epsilon}_0 \exp\left[-\frac{(t-t_m)}{\tau_m}\right] = \dot{\epsilon}_0 \exp[-\phi(t-t_m)] \quad \dots(18)$$

where  $\phi = 1/\tau_m$ ,  $t_m$  is the time at which the laser pulse has been turned off and  $\dot{\epsilon}_0$  is the strain rate at  $t = t_m$ .

Using Eq. (18), the change of pressure  $dp$  in time  $dt$  can be expressed as

$$dp = \dot{p} dt = Y \dot{\epsilon} dt = Y \dot{\epsilon}_0 \exp[-\phi(t-t_m)] dt \quad \dots(19)$$

Integrating Eq. (19) and taking  $p = p_m$  at  $t = t_m$ , we get

$$p = p_m + \frac{Y \dot{\epsilon}_0}{\phi} [1 - \exp\{-\phi(t-t_m)\}] \quad \dots(20)$$

Thus, in this case  $q = d_0 p$  and  $dq/dt = d_0 dp/dt$ , can be expressed as

$$q = d_0 p_m + \frac{d_0 Y \dot{\epsilon}_0}{\phi} [1 - \exp\{-\phi(t-t_m)\}] \quad \dots(21)$$

$$\text{and, } \frac{dq}{dt} = d_0 Y \dot{\epsilon}_0 [\exp\{-\phi(t-t_m)\}] \quad \dots(22)$$

Using Eqs. (15), (21) and (22), the EML intensity for  $q \gg q_{th}$ , can be written as

$$I = \eta \sigma \Omega n_1 n'_c n_t n_c z \tau \mu b^2 \left[ d_0 p_m + \frac{d_0 Y \dot{\epsilon}_0}{\phi} \right] [1 - \exp\{-\phi(t-t_m)\}] d_0 Y \dot{\epsilon}_0 [\exp\{-\phi(t-t_m)\}] \quad \dots(23)$$

As  $\phi$  is a very large,  $d_0 p_m \gg \frac{d_0 Y \dot{\epsilon}_0}{\phi}$ , and Eq. (23) can be written as

$$\begin{aligned} \text{or, } I &= \eta \sigma \Omega n_1 n'_c n_t n_c z \tau \mu b^2 d_0^2 p_m Y \dot{\epsilon}_0 \exp[-\phi(t-t_m)] \\ \text{or, } I &= I_m \exp[-\phi(t-t_m)] \end{aligned} \quad \dots(24)$$

where  $I_m = \eta \sigma \Omega n_1 n'_c n_t n_c z \tau \mu b^2 d_0^2 p_m Y \dot{\epsilon}_0$  is the EML intensity at  $t = t_m$ .

### EXPERIMENTAL SUPPORT TO THEORY

Fig.1 shows the mechanoluminescence glow curve induced by the application of laser impulse on  $\text{SrAl}_2\text{O}_4:\text{Eu}$ , phosphor (Akiyama et al.2002). It is seen that, initially the EML intensity increases linearly with time, attains peak value  $I_m$  for a particular time  $t_m$ . These results are in accordance with Eqs. (17), and (24).

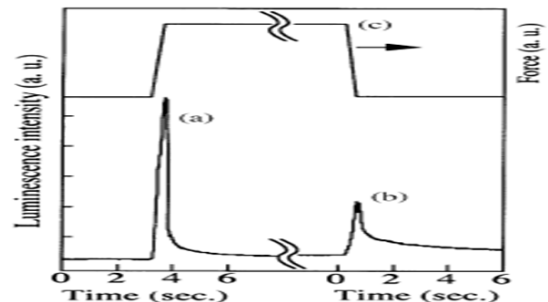


Fig 1. Mechanoluminescence glow curve of  $\text{SrAl}_2\text{O}_4:\text{Eu}$ , produced during (a) the application of laser impulse and (b) the removal of laser impulse

## CONCLUSION

The EML intensity of  $\text{SrAl}_2\text{O}_4:\text{Eu}$  phosphor excited by laser impulse, initially intensity increases with time, attains a peak value and later on it decreases with time, in which the decay of EML intensity gives the decay time of impact stress and the slow decay time of the EML intensity gives the lifetime of elections in the shallow traps lying in the normal piezoelectric region of the crystals.

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